

# Evaluation of indentation fracture toughness for brittle materials based on the cohesive zone finite element method

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## ABSTRACT

We propose indentation fracture toughness evaluation methods based on finite element analysis adapting the cohesive zone model. Establishing an appropriate finite element model, we examine the effects of material properties on the size of the cracks formed by Vickers indentation, and then suggest a regression formula for the estimation of the fracture toughness. The formula is extended to other types of pyramidal indenters including variations of the indenter angle and the number of indenter edges. The suggested methods are examined by comparison of the experimental data, and the fracture toughness can be accurately estimated from pyramidal indentation data.

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## 1. Introduction

In contrast to conventional fracture toughness evaluation methods, which require tedious preparation of the test specimens, the indentation test-based method does not require this kind of special specimen; instead, it can be directly applied to a part/structure with simple treatment of the surface. This huge advantage has led to numerous studies on evaluating the critical load corresponding to crack initiation or certain crack characteristics for Vickers, Berkovich, cube-corner, and other types of indenters. For brittle materials such as glass or ceramics, not only elastic–plastic deformation occurs during indentation due to stress concentrations beneath the indenter, but also various cracks can form around the impression. When using three- or four-sided pyramidal indenters, radial, median, semi-circular (half-penny) or lateral cracks can form independently and merge, depending on the material properties, the indenter geometry, and the indentation load. Based on these observations, numerous studies have been conducted to predict the fracture toughnesses of brittle materials from crack size, indentation load and/or other parameters [12,5,13,14,1,2,3,16,17,24,10,4,21,8]. The final goal of these studies is to find the fracture toughness similar to that obtained by conventional testing methods. However, indentation fracture toughnesses predicted by these formulas often exhibit quite a significant deviation from the fracture toughnesses of conventional methods, and therefore doubt remains regarding the reliability of the indentation cracking-based methods.

Before the 1950s, the cracks formed by indentation on brittle materials (to obtain hardness) were regarded as a disturbing phenomenon. Palmqvist [20] was the first performer who related the indentation crack length to the fracture toughness. He

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### Nomenclature

$a$	impression half-diagonal
$c$	crack length
$E$	Young's modulus
$H$	hardness
$h$	indentation depth
$K_c$	fracture toughness
$l$	crack length
$P$	indentation load
$\alpha$	correction factor
$\delta_c$	crack-initiating separation
$\delta_{\max}$	damage-initiating separation
$\varepsilon_o$	yield strain
$\Gamma$	fracture energy
$\kappa$	correction factor
$\nu$	Poisson's ratio
$\sigma_o$	yield strength
$\sigma_{\max}$	damage-initiating stress
$\psi$	centerline-to-face angle

showed that the average length of the cracks emanating from the corners of the impression can be used to infer the fracture toughness. Later, Lawn et al. [14] and Anstis et al. [1] applied Hill [6] and established a formula for calculating the fracture toughness from the crack length, the maximum indentation load, the hardness and the Young's modulus. Lawn et al. [14] findings are based on the following two experimental observations: first, the deformation under the indenter satisfies geometrical self-similarity; second, for “well-developed” cracks, the indentation load  $P$  is proportional to  $c^{3/2}$  where  $c$  is the crack length (Fig. 1). The full equation includes the Young's modulus  $E$ , the hardness  $H$  and the indenter angle  $\psi$  as follows:

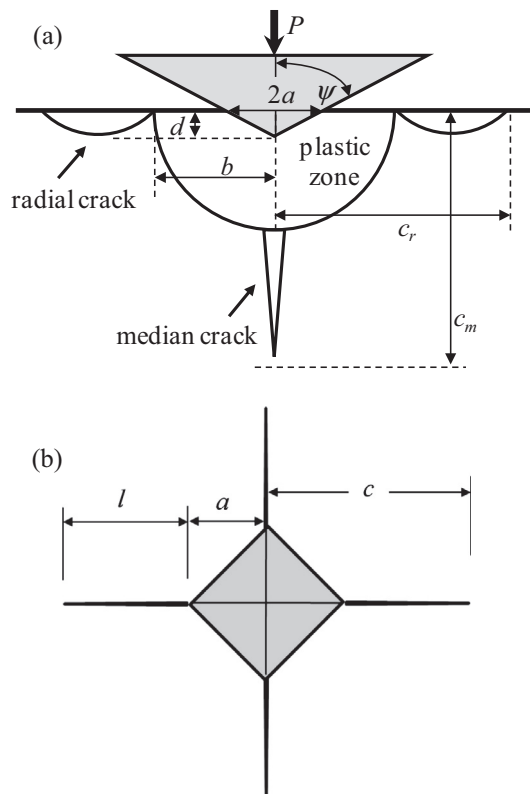


Fig. 1. Schematic of (a) the median and radial crack system and (b) the top view of Vickers indentation cracking geometry.

$$K_c \propto (a/b)(E/H) \cot \psi \frac{P}{c^{3/2}} \quad (1)$$

where  $a$  and  $b$  are the impression half-diagonal and the plastic zone size, respectively, as shown in Fig. 1. The ratio  $a/b$  is related to  $\beta_H$  ( $=b/r$  where  $r$  is the cavity radius) used in Hill's expanding cavity model, and  $\beta_H$  can be related to  $E/H$ . Based on Hill's expanding cavity model,  $H/E$  can be expressed as a function of  $b/r$ ,

$$\frac{H}{E} = \frac{2[1 + \ln(b/r)^3]}{9[(1 - \nu)(b/r)^3 - 2(1 - 2\nu)/3]} \quad (2)$$

Lawn et al. [14] simply set  $b/r \approx (E/H)^m$  and  $m = 1/2$  from the regression of experimental results. Assuming that the cavity volume is equal to the volume of the impression generated by the pyramidal indenter, we can write

$$b/a \approx (E/H)^{1/2} (\cot \psi)^{1/3}. \quad (3)$$

since  $r$  is proportional to  $a(\cot \psi/\pi)^{1/3}$ . Combining Eqs. (1) and (3) yields the following expression for the fracture toughness  $K_c$ ,

$$K_c = \alpha \left( \frac{E}{H} \right)^{1/2} \left( \frac{P}{c^{3/2}} \right). \quad (4)$$

Anstis et al. [1] calibration constant  $\alpha$  is  $0.016 \pm 0.004$  from Vickers indentation tests on glasses and ceramics. However, in contrast to Lawn et al. [14] and Anstis et al. [1], Laugier [10] used  $b/r \approx (E/H)^m$  with  $m = 1/3$ ; the exponent in Eq. (4) then becomes  $2/3$  and  $\alpha = 0.0098 \pm 0.0025$ . Eq. (4) based on Eq. (1) and Hill's expanding cavity model contains several assumptions and simplifications. Moreover, the coefficients  $m$  and  $\alpha$  were not determined by physical bases, but determined by fitting of the experimental results with restrictedly chosen materials. In each of these studies, the deviation of  $\alpha$  obtained from experimental data is about 25% although the inaccuracies in the fracture toughnesses are not merely due to  $\alpha$  but may also come from inaccurately determined material properties. In addition, Jang and Pharr [8] expressed  $\alpha$  as a function of indenter angle and Poisson's ratio, but Lee et al. [15] showed that  $\alpha$  does not only depend on Poisson's ratio but is also influenced by other material properties. Therefore, for the accurate fracture toughness evaluation, the variations of  $\alpha$  and crack morphology with respect to indenter shapes and material properties need to be thoroughly investigated. Meanwhile, Niihara et al. [16] suggested a method that distinguishes between radial ( $c/a \leq 2.5$ ) and half-penny ( $c/a \geq 2.5$ ) cracks, by using two kinds of crack lengths,  $l$  and  $c$ , as shown in Fig. 1,

$$\left( \frac{K_c \phi}{H_V a^{1/2}} \right) \left( \frac{H_V}{E \phi} \right)^{2/5} = 0.035(l/a)^{-1/2}; \quad (c/a \leq 2.5) \quad (5)$$

$$\left( \frac{K_c \phi}{H_V a^{1/2}} \right) \left( \frac{H_V}{E \phi} \right)^{2/5} = 0.129(c/a)^{-3/2}; \quad (c/a \geq 2.5) \quad (6)$$

where  $l$  is the crack length from  $a$  to the end of crack ( $c = l + a$ ),  $\phi$  ( $\approx 3$ ) is constrain factor, and  $H_V = 0.9272H$ . However, it is difficult to infer the final crack shape based on the surface cracks only because radial and median cracks can be independently initiated and propagated during loading and/or unloading, and the critical  $c/a$  value of Niihara et al. [16] is also dependent on material properties and indenter shapes. Apart from the equations above, more than 30 other forms for determining the fracture toughness by indentation tests have been proposed. Ponton and Rawlings [22,23] assessed a number of such equations and found that there is a deviation of approximately 30% from the fracture toughnesses obtained by other conventional methods.

Recently, finite element analysis (FEA) approaches and various measuring techniques for diverse indenter shapes have been applied to accurately evaluate the fracture toughness considering the effects of material characteristics [27,25,15,9]. In particular, the cohesive interface model can be efficiently applied to the FE simulation of various crack problems [26,28,15,9,7]. Lee et al. [15] have performed FE indentation cracking analysis using the cohesive zone model (CZM) to analyze the influence of the four-sided indenter angle and material properties on the crack characteristics. Johanns et al. [9] have shown that the cohesive FE model can be used to examine the relationships between crack morphology, material properties, indenter geometry, and indentation test measurements from a two-dimensional standard CT specimen, a two-dimensional wedge, and three-dimensional Vickers and Berkovich indentation cracking analyses when the crack length is long enough in comparison to the cohesive zone size. Hyun et al. [7] have extended Lee et al. [15] approach and investigated the relation between the crack length and the number of indenter edges from the CZM-based FEA and experimental comparison. With the CZM, we can simulate the fracture process in material by using the fracture energy function.

In this study, based on the model proposed by Lee et al. [15], we establish new equations to determine the fracture toughness through the quantitative evaluation of the crack size obtained by FEA with respect to material properties (the Young's modulus, the yield strength, the Poisson's ratio, and the fracture toughness) and indentation conditions (the maximum indentation load, the number of indenter edges, and the indenter angle), and then verify the equations by comparing estimated fracture toughnesses with experimental results.

The paper is organized as follows. In Section 2, the finite element analysis procedures adapting cohesive zone model are introduced. In Section 3, preliminary FE analysis is performed to find the parameters that mainly influence the fracture toughness, and their influence on the crack length is analyzed. Based on the FE parametric study of Vickers indentation, we suggest mapping functions that express the fracture toughness in terms of the maximum indentation load, the crack

length, and material properties. We study then the influence of the number of indenter edges and the angle on the crack length, and extend the applicability of the proposed equation. In Section 4, the proposed fracture toughness equations are verified by comparison to indentation test data.

## 2. Cohesive zone finite element model

Lawn et al. [14] assumed  $\alpha$  in Eq. (4) to be only dependent on the centerline-to-face angle  $\psi$ ; but as shown by Lee et al. [15], it can vary with material properties and other indentation parameters. To study the effects of material properties and indenter geometry on the fracture toughness evaluation, the commercial package, ABAQUS/Standard, is used for the cohesive zone finite element simulations. The three-dimensional FE models for four-sided and three-sided pyramidal indentation simulations are generated as shown in Fig. 2. Because of the symmetric boundary conditions, an 1/4 model for four-sided indenters and an 1/3 model for three-sided indenters are used. For the simplicity of input material properties, elastic-perfectly plastic materials are assumed. The elastic diamond indenter with the Young's modulus  $E_I = 1016$  GPa and the Poisson's ratio  $\nu_I = 0.07$  is pressed into the top surface of the specimen while the bottom surface of it is fixed.

We put the cohesive elements where the crack can be initiated and propagated. In Johanns et al. [9] and Hyun et al. [7] observation, the cracks are generated along the indenter edges only irrespective of three-sided or four-sided indenter if the symmetric conditions are maintained, so the cohesive elements are located along the plane that is perpendicular to the material free surface and aligned with the edges of the indenter, as shown in Fig. 2.

The effects of the cohesive zone properties on the crack evolution and propagation are described and discussed by Lee et al. [15] and Johanns [9]. Although the phenomenological CZM has inherent limitation for relatively short cracks in comparison to the process zone size, the variation of the traction-separation properties has a negligible effect on the indentation cracking analysis for well-developed cracks when the fracture energy and other material properties are identical. In this study, we are interested in well-developed cracks, so we simply use the bilinear traction-separation model as shown in Fig. 3. The cohesive traction increases up to a critical (damage-initiating) separation ( $\delta_{\max}$ ) between the cohesive interface, and decreases monotonically to zero upon further loading once the critical separation is exceeded. The area under the traction-separation curve in Fig. 3 represents the fracture energy  $\Gamma$ , so the fracture toughness can be inferred by

$$\Gamma = \frac{K_c^2}{E'} = \frac{1}{2} \sigma_{\max} \delta_c, \quad E' = \begin{cases} E & : \text{plane stress} \\ E/(1 - \nu^2) & : \text{plane strain} \end{cases} \quad (7)$$

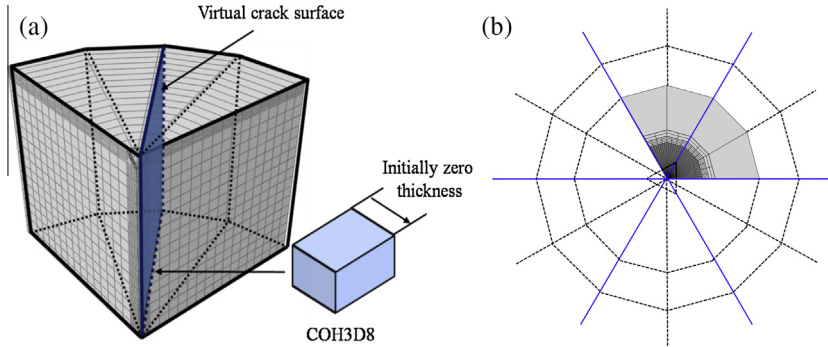


Fig. 2. (a) 1/4 FE model for four-sided and (b) 1/3 FE model for three-sided pyramidal indentation using cohesive zone model.

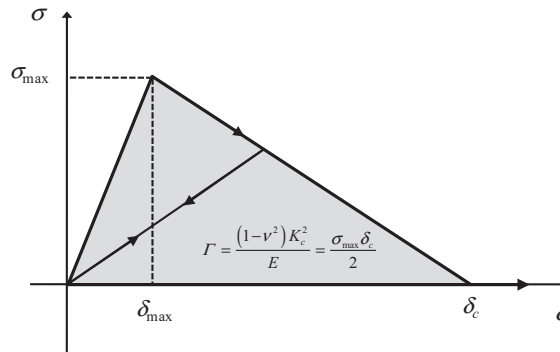


Fig. 3. The traction-separation relation for a bilinear cohesive zone model [15].

where  $\nu$  and  $E$  are the Poisson's ratio and the elastic (Young's) modulus;  $\sigma_{\max}$ ,  $\delta_{\max}$  and  $\delta_c$  are the damage-initiating stress, the corresponding separation and the crack-initiating separation, respectively, as shown in Fig. 3. The plane strain relation is used to calculate the fracture energy. It should be noted that the purpose of the indentation cracking approach to the evaluation of fracture toughness is basically to get the mode I fracture toughness value. Theoretically, the radial/median/half-penny cracks are assumed to be developed in the symmetric plane only. The mode I stress intensity factor for the center-loaded half-penny crack can be written by  $K = 2aP/(pc)^{3/2}$  which is the fundamental form of the indentation fracture toughness evaluation where  $K$  is the mode I stress intensity factor [14]. The mode-II and mode-III type cracks cannot be considered irrespective of the number of cracks [19].

Both of the 1/4 model (for four-sided indenters) and the 1/3 model (for three-sided indenters) consist of approximately 86,600 elements and 97,000 nodes. For the three-sided indenter, the crack is not developed on the face plane, but developed on the edge plane only when the symmetric conditions are assumed in FEA, so we do not consider the crack development on the face plane, and the 1/3 model is used in the present work [7].

The detailed procedure to measure the crack length in FEA is also described in Lee et al. [15] and Hyun et al. [7] papers. The crack length depends on the definition of the crack front around the process zone in FEA, but the sensitivity to the definition decreases with the process zone size. In this study, we are interested in well-developed cracks, so the definition is not a critical factor. We simply assume that the crack forms when the cohesive traction has become zero, i.e.,  $\delta = \delta_c$ . The damage-initiating stress is chosen as  $\sim 0.1\sigma_o \leq \sigma_{\max} \leq \sim 0.2\sigma_o$  ( $\sigma_o$ : yield strength) to ensure the crack generation and to satisfy the linear elastic fracture condition, i.e.,  $c \gg$  process zone size [15].

The contact area depends on the friction coefficient between the indenter and the specimen, and as a consequence, the stress state in the material varies with the friction coefficient, which again influences the crack size. As the Coulomb friction coefficient  $f$  increases, the crack length very slightly increases in FE analysis; however, when  $f \geq 0.2$ , the crack length does not increase further. The friction coefficient between the diamond and other materials is about 0.1, so  $f$  is assumed to be 0.1.

### 3. Evaluation of indentation fracture toughness based on FEA

#### 3.1. Preliminary FE analysis for fracture toughness evaluation

Preliminary FE analysis is performed to set up indentation parameters for fracture toughness evaluation. We set the damage-initiating stress and the fracture energy as  $\sigma_{\max} = 0.5$  GPa and  $\Gamma = 0.0025$  GPa  $\mu\text{m}$ , respectively, based on Lee et al. [15] work. When the fracture toughness, the yield strain ( $\varepsilon_o \equiv \sigma_o/E$ ), the Poisson's ratio, and the indenter centerline-to-face angle are changed, we investigate the variation of the coefficient of Eq. (4),  $\alpha$ .

When the Young's moduli ( $E$ ) are 100, 200, 400 GPa and the Poisson's ratio is fixed as 0.3, the fracture toughness ( $K_{IC}$ ) values become 0.52, 0.74, 1.05 MPa  $\text{m}^{1/2}$ , respectively. Table 1 shows the effects of the fracture toughness and the yield strain on the coefficient  $\alpha$ . Other parameters used in the simulation are also listed in Table 1. For fixed  $\varepsilon_o$  and  $\Gamma$ ,  $\alpha$  decreases with increasing  $E$ , and for fixed  $E$  and  $\Gamma$ ,  $\alpha$  decreases with increasing  $\varepsilon_o$  (i.e.,  $\sigma_o$ ). This means not only  $\varepsilon_o$  but also  $E$  itself has an influence on  $\alpha$ . Table 2 shows when other material properties and parameters are fixed,  $\alpha$  decreases with increasing the Poisson's ratio  $\nu$ . As mentioned by Lee et al. [15], the crack morphology is significantly dependent on the Poisson's ratio. Hence, the Poisson's ratio is another important parameter to accurately evaluate the fracture toughness whereas the value is simply

**Table 1**

The variations of  $c/a$  and  $\alpha$  with respect to material properties  $K_{IC}$  and  $\varepsilon_o$  for Vickers indentation.

$\nu = 0.3$ , $\Gamma = 0.0025$ GPa $\mu\text{m}$ $h_{\max} = 0.6$ $\mu\text{m}$ , $\sigma_{\max} = 0.5$ GPa, $\delta_{\max}/\delta_c = 1/4$					
$K_{IC}$ (MPa $\text{m}^{1/2}$ )	$\varepsilon_o$ ( $\sigma_o/E$ )	$P_{\max}$ (mN)	$H$ (GPa)	$c/a$	$\alpha$
0.52	0.025 (2.5/100)	42.6	5.7	1.58	0.0157
0.74	0.025 (5.0/200)	84.8	11.3	1.92	0.0149
1.05	0.025 (10/400)	168.8	22.5	2.23	0.0132
0.52	0.050 (5.0/100)	55.0	8.2	1.53	0.0127
0.74	0.050 (10/200)	113.9	17.0	1.83	0.0116
1.05	0.050 (20/400)	227.4	34.0	2.11	0.0102

**Table 2**

The variations of  $c/a$  and  $\alpha$  with respect to Poisson's ratio for Vickers indentation.

$E = 200$ GPa, $\sigma_o = 5$ GPa, $\Gamma = 0.0025$ GPa $\mu\text{m}$ $h_{\max} = 0.6$ $\mu\text{m}$ , $\sigma_{\max} = 0.5$ GPa, $\delta_{\max}/\delta_c = 1/4$					
$K_{IC}$ (MPa $\text{m}^{1/2}$ )	$\nu$	$P_{\max}$ (mN)	$H$ (GPa)	$c/a$	$\alpha$
0.71	0.1	76.3	10.2	2.25	0.0191
0.72	0.2	80.1	10.7	2.10	0.0170
0.74	0.3	84.8	11.3	1.93	0.0150

**Table 3**The variations of  $c/a$  and  $\alpha$  with respect to Vickers indenter angle.

$E = 200 \text{ GPa}$ , $\nu = 0.3$ , $\sigma_0 = 5 \text{ GPa}$ , $K_c = 0.74 \text{ MPa m}^{1/2}$ ( $\Gamma = 0.0025 \text{ GPa } \mu\text{m}$ ) $h_{\max} = 1.0 \text{ } \mu\text{m}$ , $\sigma_{\max} = 0.5 \text{ GPa}$ , $\delta_{\max}/\delta_c = 1/4$				
$\psi$ ( $^\circ$ )	$P_{\max}$ (mN)	$H$ (GPa)	$c/a$	$\alpha$
55	99.2	11.8	2.72	0.0217
68	231.7	10.6	2.45	0.0159
75	426.1	10.3	2.17	0.0107

assumed to be 0.25 in the LEM form, Eq. (4). Table 3 shows the variation of  $\alpha$  with respect to the centerline-to-face angle  $\psi$ . As the indenter angle increases,  $\alpha$  decreases. This is because at the same maximum indentation loads, lower  $\psi$  results in a larger indent volume, which produces longer cracks and a larger  $\alpha$ .

Consequently, for the accurate evaluation of the fracture toughness, the evaluation formulas should be set up to take into account the effects of the yield strain  $\varepsilon_0$ , the Young's modulus  $E$ , the Poisson's ratio  $\nu$ , and indenter angle  $\psi$ .

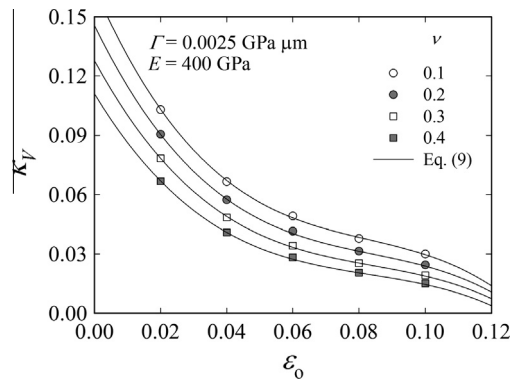
### 3.2. Fracture toughness evaluation forms based on a parametric study

In Section 3.1, we have shown that  $\alpha$  in Eq. (4) is a function of  $\varepsilon_0$ ,  $E$ ,  $\nu$  and  $\psi$ . However, the hardness in Eq. (4) is not a material property because it can vary with indenter geometry, so it is a function of other material properties. In the present work, we preferentially reform the evaluation formula excluding the hardness. Based on the relationship between the maximum indentation load and the crack length for the well-developed crack system, i.e.,  $P_{\max} \propto c^{3/2}$  [11], we can express the fracture toughness as a function of  $\cot \psi$ ,  $\nu$ ,  $\varepsilon_0$ , and  $E$  as follows:

$$\frac{K_c}{P_{\max}/c^{3/2}} = \kappa(\cot \psi, \nu, \varepsilon_0, E). \quad (8)$$

In this section, we investigate the effect of each parameter on the coefficient  $\kappa$ , and then establish a new relation to evaluate the fracture toughness based on the FE parametric study. Afterwards, we suggest another equation to relate the material properties to the hardness.

FE analyses for four-sided pyramidal indentation are performed for each combination of five yield strains ( $\varepsilon_0 = 0.02, 0.04, 0.06, 0.08, 0.1$ ), four Poisson's ratios ( $\nu = 0.1, 0.2, 0.3, 0.4$ ), two centerline-to-face angles ( $\psi = 55^\circ, 68^\circ$ ), and four Young's moduli ( $E = 100, 200, 400, 600 \text{ GPa}$ ). Jang and Pharr [8] experimental and Hyun et al. [7] FE results confirmed that the crack length  $c$  is proportional to  $(\cot \psi)^{4/9}$ , so the coefficient  $\kappa$  in Eq. (8) is proportional to  $(\cot \psi)^{2/3}$ . We therefore first establish an equation to predict fracture toughness for Vickers indenters ( $\psi = 68^\circ$ ), and then the equation is extended to other types of indenters by using  $\kappa \propto (\cot \psi)^{2/3}$ . We set  $\Gamma = 0.0025, 0.0010 \text{ GPa } \mu\text{m}$  and  $E = 100, 200, 400, 600 \text{ GPa}$ , so the fracture toughness becomes around  $0.5\text{--}2 \text{ MPa } \mu\text{m}^{1/2}$ , which depends on the Poisson's ratio. The materials are indented until the cracks are well-developed, and  $\kappa$  values are then evaluated. For each  $E$ , we observe the influences of  $\varepsilon_0$  and  $\nu$  on  $\kappa_V$  where  $\kappa_V$  is the  $\kappa$  obtained with Vickers indentation (centerline-to-face angle  $\psi = 68^\circ$ ). Fig. 4 shows  $\kappa_V\text{--}\varepsilon_0$  relation with respect to the Poisson's ratio for  $E = 400 \text{ GPa}$ . When  $\Gamma$ ,  $E$  and  $\varepsilon_0$  are fixed,  $\kappa_V$  decreases with increasing  $\nu$ . This can be explained by the fact that the decrease in  $\nu$  makes cracks further developed during the unloading as described by Lee et al. [15], and because the load deviation caused by the change of the Poisson's ratio is relatively small,  $P_{\max}/c^{3/2}$  increases with  $\nu$  which makes  $\kappa_V$  decreased. When  $\Gamma$ ,  $E$  and  $\nu$  are fixed,  $\kappa_V$  decreases with increasing  $\varepsilon_0$ , which means the increase in the yield resistance results in shorter cracks.



**Fig. 4.**  $\kappa_V$  vs.  $\varepsilon_0$  with respect to Poisson's ratio for Vickers indentation ( $\Gamma = 0.0025 \text{ GPa } \mu\text{m}$ ,  $E = 400 \text{ GPa}$ ). The regression lines obtained from Eq. (9) are also plotted.

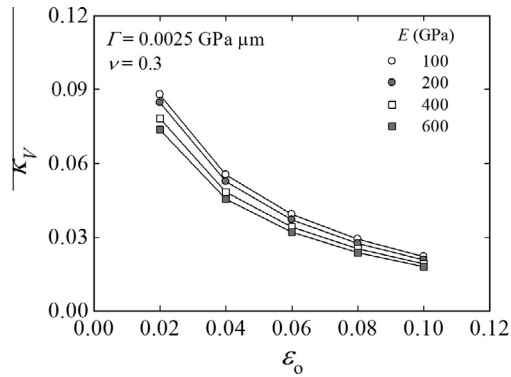


Fig. 5.  $\kappa_V$  vs.  $\varepsilon_o$  with respect to Young's modulus for Vickers indentation ( $\Gamma = 0.0025$  GPa  $\mu\text{m}$ ,  $\nu = 0.3$ ).

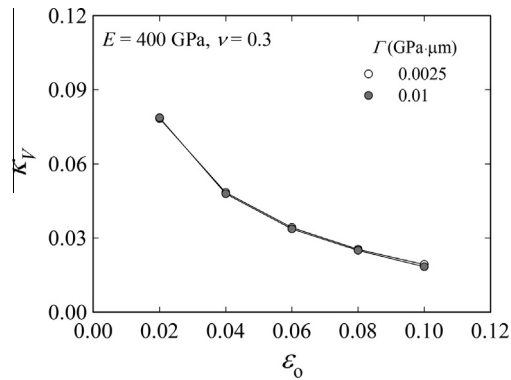


Fig. 6.  $\kappa_V$  vs.  $\varepsilon_o$  with respect to fracture toughness for Vickers indentation ( $E = 400$  GPa,  $\nu = 0.3$ ).

When the material properties are fixed as  $\Gamma = 0.0025$  GPa  $\mu\text{m}$  and  $\nu = 0.3$ , the variations of  $\kappa_V$  with respect to  $E$  and  $\varepsilon_o$  are examined. As shown in Fig. 5, when  $\varepsilon_o$  is fixed,  $\kappa_V$  decreases with increasing  $E$ , which is the same as the results from the LEM form in Section 3.1. Fig. 6 shows the effect of the  $\Gamma$  (or  $K_c$ ) on  $\kappa_V$ . When  $E = 400$  GPa and  $\nu = 0.3$ ,  $\kappa_V$  values are obtained for the five yield strains. The  $\kappa_V$  values obtained with  $\Gamma = 0.010$  GPa  $\mu\text{m}$  ( $K_c = 2.0966$  MPa  $\text{m}^{1/2}$ ) are in good agreement with them with  $\Gamma = 0.0025$  GPa  $\mu\text{m}$  ( $K_c = 1.0482$  MPa  $\text{m}^{1/2}$ ). Thus, when  $E$ ,  $\nu$  and  $\varepsilon_o$  are fixed,  $\kappa_V$  values are independent of  $\Gamma$  (or  $K_c$ ), which means Eq. (8) is a reasonable form. To consider the influence of  $E$  on  $\kappa_V$ , we use  $E_R$ , the ratio of  $E$  to a reference Young's modulus, 1000 GPa ( $E_R = E/E_{1000}$ ), which is similar to the indenter modulus, for the normalization of  $E$ .

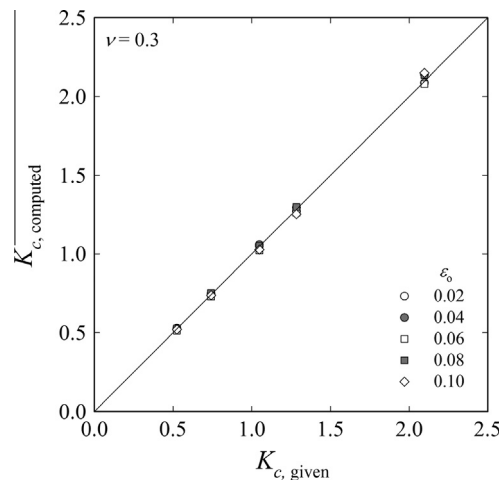


Fig. 7. The comparison of the fracture toughnesses computed from Eq. (9) to those imposed on FE model.



### 3.3. Regression analyses

The following procedure is used for the regression of each variable. First, for a given yield strain and Poisson's ratio, we regress to  $E_R$ . The obtained regression coefficients are regressed to the yield strain, and finally to the Poisson's ratio. Eq. (9) is a unified equation that takes into account each variable, i.e., the Young's modulus ratio, the yield strain and the Poisson's ratio.

$$\begin{aligned} K_c &= \kappa_V(v, \varepsilon_0, E_R) P_{\max} / c^{3/2} \\ \kappa_V(v, \varepsilon_0, E_R) &= \alpha_i^K(v, \varepsilon_0) (E_R)^i; \quad i = 0, 1, 2 \\ \alpha_i^K(v, \varepsilon_0) &= \beta_{ij}^K(v) \varepsilon_0^j; \quad j = 0, 1, 2, 3 \\ \beta_{ij}^K(v) &= \gamma_{ijk}^K v^k; \quad k = 0, 1, 2 \end{aligned} \quad (9)$$

where we use the Einstein summation convention, and  $\alpha_i$ ,  $\beta_{ij}$ , and  $\gamma_{ijk}$  are coefficients of the polynomial. The values of coefficients are given in Appendix A. The regression lines for  $\Gamma = 0.0025$  GPa  $\mu\text{m}$  and  $E = 400$  GPa are depicted in Fig. 4, and the input values of FEA and predicted fracture toughnesses are plotted in Fig. 7. The average deviation between calculated  $K_c$  values and those given is 1.2%.

As described before,  $\kappa$  is linear to  $(\cot \psi)^{2/3}$ . Hence, for example, the relationship between the  $\kappa$  values obtained with the  $\psi = 55^\circ$  four-sided indenter and the  $\psi = 68^\circ$  Vickers indenter can be derived:

$$\kappa_{4-55^\circ} = \kappa_V \left( \frac{\cot 55^\circ}{\cot 68^\circ} \right)^{2/3}. \quad (10)$$

where the subscript 4–55° means the 55° four-sided indenter. Hence, Eq. (9) based on Vickers indentation becomes

$$K_c = \kappa_V(v, \varepsilon_0, E_R) \left( \frac{\cot 55^\circ}{\cot 68^\circ} \right)^{2/3} P_{\max} / c^{3/2}. \quad (11)$$

By using the maximum load and the crack length obtained from FE simulation with  $\psi = 55^\circ$  four-sided pyramidal indenter, the fracture toughnesses are calculated as listed in Table 4. The average and the maximum deviations between input and predicted fracture toughnesses are about 2% and 4%, respectively.

Another factor related to the indenter geometries is the number of the edges (cracks),  $n_c$ . Hyun et al. [7] have suggested the following equation,

$$c_{Vn_c} / c_V = k_1 e^{-k_2 n_c}; \quad (k_1, k_2) = (1.3680, 0.0778), \quad n_c \geq 3 \quad (n_c \text{ is an integer}) \quad (12)$$

where  $c_V$  is the crack length obtained from Vickers indentation, and  $c_{Vn_c}$  is the crack length that can be obtained from another pyramidal indenter having  $n_c$  edges when their ideal contact areas (without consideration of the pile-up/sink-in effect) are equivalent. For example,  $n_c = 3$  means the Berkovich indenter. When applying the relation between the crack size and the number of cracks at the same indentation load, we obtain the following fracture toughness equation for Berkovich indentation;

$$K_c = \kappa_B(v, \varepsilon_0, E_R) P_{\max} / c_B^{3/2} = \kappa_V(v, \varepsilon_0, E_R) (k_1 e^{-3k_2})^{3/2} P_{\max} / c_B^{3/2} \quad k_1, k_2 = 1.3680, 0.0778 \quad (13)$$

where the subscript  $B$  means a Berkovich indenter. If an  $n_c$ -sided indenter with arbitrary centerline-to-face angle  $\psi$  is used, Eq. (13) can be further revised as follows:

$$K_c = \kappa_{n_c}(v, \varepsilon_0, E_R) P_{\max} / c_{n_c}^{3/2} = \kappa_V(v, \varepsilon_0, E_R) (k_1 e^{-n_c k_2})^{3/2} \left( \frac{\cot \psi_{n_c}}{\cot \psi_{Vn_c}} \right)^{2/3} P_{\max} / c_{n_c}^{3/2} \quad k_1, k_2 = 1.3680, 0.0778 \quad (14)$$

where subscript  $n_c$  means an  $n_c$ -sided indenter, and  $Vn_c$  means an  $n_c$ -sided indenter equivalent to the Vickers indenter.

**Table 4**

The comparison of the fracture toughnesses obtained from Eq. (11) to those given for  $\psi = 55^\circ$  four-sided pyramidal indentation.

$v = 0.3$ , $\varepsilon_0 = 0.04$ , $\Gamma = 0.0025$ GPa $\mu\text{m}$ , $\sigma_{\max} = 0.5$ GPa, $\delta_{\max} / \delta_c = 1/4$					
Given $K_c$ (MPa $\text{m}^{1/2}$ )	$h_{\max}$	$E$ (GPa)	$\kappa_{4-55^\circ}$	Computed $K_c$ (MPa $\text{m}^{1/2}$ )	Error (%)
0.5241	1.5	100	0.0831	0.5148	1.8
0.7412		200	0.0780	0.7297	1.6
1.0483		400	0.0718	1.030	1.8
1.2839		600	0.0691	1.233	4.0
2.0966 <sup>a</sup>		400	0.0712	2.077	0.94

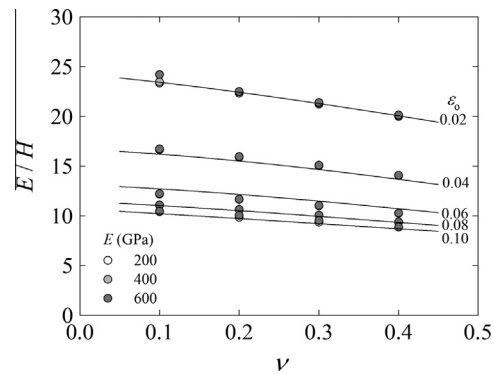
<sup>a</sup>  $\Gamma = 0.010$  GPa  $\mu\text{m}$ .



**Table 5**

The comparison of fracture toughnesses obtained from Eq. (13) to those given for Berkovich indentation.

$\nu = 0.3, \varepsilon_0 = 0.04, I^* = 0.0025 \text{ GPa} \cdot \mu\text{m}, \sigma_{\max} = 0.5 \text{ GPa}, \delta_{\max}/\delta_c = 1/4$					
Given $K_c$ (MPa m <sup>1/2</sup> )	$h_{\max}$	$E$ (GPa)	$\kappa_B$	Computed $K_c$ (MPa m <sup>1/2</sup> )	Error (%)
0.5241	1.5	100	0.0657	0.5028	4.1
0.7412		200	0.0614	0.7242	2.3
1.0483		400	0.0575	1.005	4.1
1.2839		600	0.0542	1.226	4.5

**Fig. 8.**  $E/H$  vs.  $\nu$  for five yield strains. The effect of Young's modulus variation is negligible. The regression lines obtained from Eq. (15) are also plotted.

The maximum load and the crack length are acquired from the Berkovich indentation simulations, which are substituted into Eq. (13) to predict the fracture toughness. As shown in Table 5, the maximum deviation between input and calculated fracture toughnesses is below 5%.

Eqs. (9), (13) and (14) use the yield strain  $\varepsilon_0$  as an input parameter, but it is an additional parameter that cannot be directly obtained from the indentation test. We therefore replace the regression parameter  $\varepsilon_0$  by  $E/H$  for Vickers indentation, thereby making the evaluation method more practical. The relation between  $E/H$ , the Poisson's ratio and the yield strain is observed by FE analysis. From Fig. 8, it can be seen that  $E/H$  decreases with increasing  $\nu$  and  $\varepsilon_0$ , and when  $\varepsilon_0$  is fixed,  $E/H$  is almost constant irrespective of  $E$ . For four Poisson's ratios and five yield strains, we establish a relation between material properties and  $E/H$  using a regression method similar to previous equations. The regression formula becomes

$$\begin{aligned} E/H &= f_0(\nu) + f_1(\nu)e^{-f_2(\nu)\varepsilon_0} \\ f_i(\nu) &= g_{ij}\nu^j; \quad i, j = 0, 1, 2 \end{aligned} \quad (15)$$

where we use the Einstein summation convention, and  $g_{ij}$  is a coefficient of the polynomial. The values of coefficients are given in Appendix A. Eq. (15) can be reversely written as follows:

$$\begin{aligned} \varepsilon_0 &= -\frac{1}{f_2(\nu)} \ln \left( \frac{E/H - f_0(\nu)}{f_1(\nu)} \right) \\ f_i(\nu) &= g_{ij}\nu^j; \quad i, j = 0, 1, 2. \end{aligned} \quad (16)$$

Eq. (16) enables us to predict the fracture toughness from  $E/H$ , which can be directly obtained from the Vickers indentation test. It should be noted that we basically assume elastic-perfectly plastic materials in FE simulation. Hence, the yield strain

**Table 6**

The comparison of the fracture toughnesses obtained by Eq. (9) to those given in Anstis et al. [1] for Vickers indentation.

Material	$H^a$	$\sigma_o^b/E^a$	$\nu^c$	$P_{\max}/c^{3/2a,c}$	$K_c$ (MPa m <sup>1/2</sup> )	
					Traditional method <sup>a</sup>	Eq. (9)
Si <sub>3</sub> N <sub>4</sub> (NC132)	18.5	10.2/300	0.24	60	4.0	3.83 (−4.2)
Glass–ceramic (C9606)	8.4	4.5/108	0.24	43	2.5	2.52 (+1.0)
Si <sub>3</sub> N <sub>4</sub> (NC350)	9.6	4.5/170	0.24	31	2.0	2.47 (+24)
Aluminosilicate glass	6.6	3.6/89	0.28	19	0.91	1.09 (+20)
Soda-lime glass	5.5	3.1/70	0.24	14	0.74	0.79 (+7.3)

<sup>a</sup> Anstis et al. [1] ( $P_{\max}/c^{3/2}$  values are estimated from the corresponding figure in the paper).

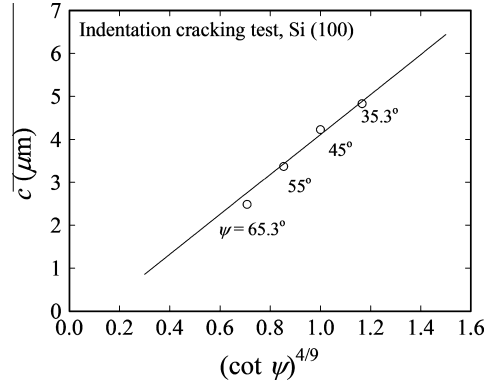
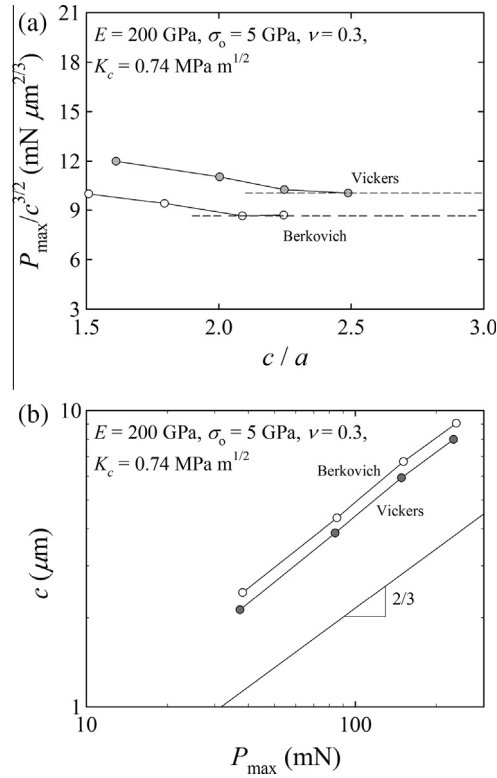
<sup>b</sup> The values obtained by trial and error method using FE analysis in the present work.

<sup>c</sup> Lee et al. [15].

**Table 7**

The comparison of fracture toughnesses obtained by Eq. (14) to a reference value for three-sided pyramidal indentation.

Indentation cracking test, Si (100), $K_c = 0.7 \text{ MPa m}^{1/2}$						
$\psi$ ( $^\circ$ )	$P_{\max}$ (mN)	$c$ ( $\mu\text{m}$ )	$c/a$	$\kappa_3$	Computed $K_c$ ( $\text{MPa m}^{1/2}$ )	Error (%)
35.3	50	4.83	2.8	0.1495	0.704	0.6
45		4.23	2.4	0.1188	0.683	−2.4
55		3.37	1.9	0.0937	0.757	8.1
65.3		2.48	1.4	0.7078	0.906	29.5


**Fig. 9.** The variation of  $c$  with respect to centerline-to-face angle  $\psi$ .

**Fig. 10.** (a)  $P_{\max}/c^{3/2}$  vs.  $c/a$  and (b)  $c$  vs.  $P_{\max}$  for Vickers and Berkovich indentation.

calculated from Eq. (16) is an apparent yield strain converted from the hardness value, so the actual yield strain is always smaller than that calculated from Eq. (16). The usage of Eq. (16) is discussed in Section 4.

**Table 8**

The comparison of fracture toughnesses obtained from the combination of Eqs. (9) and (16) to those given in Anstis et al. [1] for Vickers indentation.

Material	$P_{\max}/c^{3/2a}$	$E/H^b$	$\epsilon_0^b$	$\nu^c$	$K_c$ (MPa m <sup>1/2</sup> )		
						Traditional method <sup>a</sup>	Eq. (9)
Si <sub>3</sub> N <sub>4</sub> (NC132)	60	16.2	0.036	0.24	4	3.67 (−8.2)	3.87 (−3.4)
SiC	50	18.2	0.032	0.14	4	3.69 (−7.8)	3.41 (−15)
Al <sub>2</sub> O <sub>3</sub> (AD999)	34	20.2	0.025	0.22	3.9	2.68 (−31)	2.44 (−37)
Al <sub>2</sub> O <sub>3</sub> (AD90)	30	29.8	0.0075	0.22	2.9	3.57 (+23)	2.62 (−9.7)
Glass ceramic	43	12.8	0.053	0.24	2.5	2.05 (−18)	2.47 (−1.3)
Si <sub>3</sub> N <sub>4</sub> (NC350)	31	17.7	0.031	0.24	2	2.24 (+12)	2.09 (+4.4)
Sapphire	21	19.5	0.026	0.22	2.1	1.58 (−25)	1.48 (−29)
Aluminosilicate glass	19	13.5	0.047	0.28	0.91	0.96 (+5.4)	1.12 (+23)
Soda-lime glass	14	12.7	0.054	0.24	0.74	0.67 (−9.3)	0.80 (+8.0)
Lead alkali glass	14	13.3	0.050	0.24	0.68	0.72 (+5.2)	0.82 (+20)
Si	12	15.8	0.038	0.22	0.7	0.76 (+8.9)	0.76 (+9.2)

<sup>a</sup> Anstis et al. [1] ( $P_{\max}/c^{3/2}$  values are estimated from the corresponding figure in the paper).

<sup>b</sup> The values obtained by Eq. (16) in the present work.

<sup>c</sup> Lee et al. [15].

#### 4. Evaluation of the new formulas by using indentation test data

We compare the suggested equations to the experimental data of Anstis et al. [1] and our nanoindentation data. We first investigate the usage of Eq. (9) by comparison of Anstis et al. Vickers indentation results for brittle materials. In their Vickers indentation cracking tests, the fracture toughnesses are determined from the maximum indentation load, the crack lengths, the Young's modulus, and the hardness. However, to use Eq. (9), the yield strain and Poisson's ratio should be known values. We therefore guess the apparent yield strength (and the yield strain) that gives the hardness  $H$  from Vickers indentation FE analysis, and choose the Poisson's ratios which Lee et al. [15] have used as listed in Table 6. Anstis et al. tested fifteen materials, but it is tedious work to determine the yield strength by a trial and error method using FE analysis, so we choose only five materials and compare the results. The fracture toughnesses of the five materials that Anstis et al. [1] tested [Si<sub>3</sub>N<sub>4</sub> (NC132), Glass-ceramic (C9606), Si<sub>3</sub>N<sub>4</sub> (NC350), Aluminosilicate glass, Soda-lime glass] are in the range 0.74–4.0 MPa μm<sup>1/2</sup>, and the  $H/E$  ratios lies between 0.6 and 0.8. The average and standard deviations between the fracture toughnesses obtained from traditional methods and Eq. (9) are about 9.5% and 11.9%, respectively. Except for Si<sub>3</sub>N<sub>4</sub> (NC350), Eq. (9) yields relatively accurate fracture toughnesses. The error in fracture toughness may stem from errors in measurement of Poisson's ratio, Young's modulus and from the conventionally obtained fracture toughness itself as well as from the measured crack length.

We investigate then the effectiveness of Eq. (14) for three-sided pyramidal indenters extended from the Vickers indentation formula. Based on the nanoindentation test data on (100) Si ( $K_{c, \text{conventional}} = 0.7$  MPa m<sup>1/2</sup>) using three-sided indenters with various indenter angles, we predict the fracture toughness. The Nanoindenter XP (Agilent Technologies) is employed for the nanoindentation tests. Indenter angles are  $\psi = 35.3^\circ$  (cube-corner),  $45^\circ$ ,  $55^\circ$ ,  $65.3^\circ$  (Berkovich). The crack lengths are measured after indentation with  $P_{\max} = 50$  mN and subsequent load release. Total nine indentation tests are conducted for each indenter angle with loading rate  $\dot{\nu}_i = 0.5$  mN/s. The measured  $c$  and  $a$ , and fracture toughness values obtained by Eq. (14) are listed in Table 7. Material properties are taken from the values obtained by Hyun et al. [7] using the Oliver–Pharr method [18] and a trial-and-error method. ( $E = 138$  GPa,  $\sigma_0 = 5.4$  GPa, and  $\nu = 0.22$ ) Except for the fracture toughness calculated with  $\psi = 65.3^\circ$ , the deviation does not exceed 10% for other indenters. As shown in Table 7, the  $c/a$  ratio decreases with increasing indenter angle, and the ratio for  $\psi = 65.3^\circ$  is 1.4 which is much smaller than 2.5. Hence, in the indentation tests with  $\psi = 65.3^\circ$  when  $P_{\max} = 50$  mN, the cracks would not be well-developed, so that the correlation of  $c \propto (\cot \psi)^{4/9}$  could not be established in Fig. 9. For this reason, the prediction error becomes greater as the indenter angle increases. Jang and Pharr [8] have mentioned that even though the  $c/a$  ratio is much smaller than that required to be “well-developed” ( $c/a \ll 2.5$ ), the proportionality between  $P$  and  $c^{3/2}$  is maintained and the fracture toughness can be therefore estimated. However, Lee et al. [15] and Hyun et al. [7] have analyzed the critical  $c/a$  ratio where the maximum indentation load  $P_{\max}$  is proportional to  $c^{3/2}$  as shown in Fig. 10(a), and have found that the estimated fracture toughness significantly depends on the  $c/a$  ratio if the crack is not well-developed. This discrepancy between Jang and Pharr's and our results is caused by the difference of the plotting method. Whereas linear scales are used on both the horizontal and vertical axes in Fig. 10(a), log-log scales are used in Fig. 10(b). In the log-log plot, the dependency on the  $c/a$  ratio seems to be negligible although the deviation of the  $P_{\max}/c^{3/2}$  in Fig. 10(a) is about 20%. Hence, to achieve accurate fracture toughness values, it is necessary to use saturated data in the linear plot. For Si, the  $c/a$  ratio and the maximum load should be greater than 2 and 150–200 mN, respectively. Consequently, for well-developed cracks only, Eq. (14) can be applied to the calculation of the fracture toughness from indentation tests for various pyramidal indenter angles and edge numbers.

In Eqs. (15) and (16), we have suggested another method to estimate the yield strain from the hardness value, which can make the suggested method more practical. To verify the usefulness of the suggested method, we calculate the fracture toughnesses from Anstis et al. [1] experimental data. The used values of  $P_{\max}/c^{3/2}$ ,  $E/H$ ,  $\nu$ ,  $K_c$  and the apparent yield strains and fracture

**Table A.1**  
Coefficients of Eq. (9).

	$k = 0$	$k = 1$	$k = 2$
$i = 0$			
$j = 0$	2.093e−1	−2.119e−1	−2.500e−3
$j = 1$	−4.829e+0	3.939e+0	3.525e+0
$j = 2$	5.412e+1	−5.001e+1	−5.050e+1
$j = 3$	−2.243e+2	2.278e+2	2.342e+2
$i = 1$			
$j = 0$	−1.332e−01	3.993e−01	−3.425e−01
$j = 1$	4.647e+00	−1.568e+01	3.740e+00
$j = 2$	−6.397e+01	1.963e+02	6.472e+01
$j = 3$	2.860e+02	−7.237e+02	−8.320e+02
$i = 2$			
$j = 0$	1.768e−01	−9.491e−01	1.195e+00
$j = 1$	−8.048e+00	4.249e+01	−4.250e+01
$j = 2$	1.175e+02	−5.650e+02	4.388e+02
$j = 3$	−5.358e+02	2.303e+03	−1.204e+03

**Table A.2**  
Coefficients of Eq. (15).

	$j = 0$	$j = 1$	$j = 2$
$i = 0$	9.990e+0	−5.538e+0	2.325e+0
$i = 1$	3.030e+1	−1.532e+1	9.000e+0
$i = 2$	3.767e+1	−1.877e+1	5.075e+1

toughnesses predicted by the combination of Eqs. (16) and (9) are listed in Table 8. The average and standard deviations of suggested method are 4% and 16.6% whereas the values of the LEM method are 2.9% and 18.9%. The deviation may stem from the estimated yield strain obtained by Eq. (16) because the effect of the material hardening is neglected. The errors can also be caused by the other data, such as the Poisson's ratio, the Young's modulus, the hardness, the fracture toughness, and  $P_{\max}/c^{3/2}$ . However, it should be noted that the standard deviation of our method is smaller than that of the LEM method, even though our prefactor  $\kappa$  is purely obtained from FEA whereas the LEM prefactor  $\alpha$  is an average value of their experimental results. Consequently, considering the errors of other material properties including fracture toughness measured by conventional methods, Eqs. (9) and (16) are very accurate and useful equations to estimate the fracture toughnesses.

The unified Eq. (14) that can be applied to all kinds of pyramidal indenters is defined by only material properties and indenter geometries. Eqs. (15) and (16), however, are valid for Vickers indentation only because the material hardness is not a material property. It is varied with the indenter angle and the number of indenter edges, so it is therefore required to individually set up the form to replace  $\varepsilon_0$  by  $E/H$  for other types of pyramidal indenters.

## 5. Concluding remarks

In this study, we presented indentation fracture toughness evaluation methods based on FE analysis using the cohesive zone model. We studied the influence of material properties and indenter geometries on the crack length, and set up indentation variables. We suggested some kinds of formulas to estimate the fracture toughness. The fracture toughnesses were then estimated from the indentation test data, and compared to conventional values. We drew the following conclusions:

- The crack size is influenced by yield strain  $\varepsilon_0$ , Poisson's ratio  $\nu$ , and Young's modulus  $E$ ;  $\varepsilon_0$ ,  $\nu$ ,  $E_R$  ( $\equiv E/E_{1000}$ ) were chosen as main parameters, and fracture toughness formulas were established. In experimental data yield strain  $\varepsilon_0$  is unknown, so  $\varepsilon_0$  was determined by comparison of material hardness. The average and standard deviations between the fracture toughnesses obtained from traditional testing methods and Eq. (9) is about 10% and 12%, respectively.
- Using the relation between the number of cracks and the crack size for the equivalent indenter that gives the same ideal contact area as well as the relation between the indenter angle and the crack size for the indenters having the same number of indenter edges, the suggested fracture toughness evaluation methods for Vickers indentation can be extended to any kinds of pyramidal indenters. The unified form Eq. (14) is very accurate for well-developed cracks.
- Eqs. (9) and (14) require  $\varepsilon_0$  as an input value. However, it is difficult to determine  $\varepsilon_0$  for brittle materials, and it means additional experimental data is needed. We therefore proposed a method to derive  $\varepsilon_0$  from  $E/H$  and  $\nu$ . By combination of Eqs. (9) and (16) and Anstis et al.'s Vickers indentation test data, we evaluated and compared the fracture toughnesses; the average and standard deviation is approximately 4% and 17%, respectively. Hence, without experimental calibration of the prefactor  $\kappa$ , we can accurately evaluate the fracture toughness of brittle materials.

- (iv) In this study, we simply assumed isotropic elastic-perfectly plastic materials. The material hardening and anisotropy would affect the fracture behavior, so they would be our future work.

## Acknowledgements

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## Appendix A

The coefficient values of regression functions of Eqs. (9) and (15) (see Tables A.1 and A.2).

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